

Fully Tunable LiNbO₃ Ring Resonator Cavity for Frequency Comb Generator (FCG)

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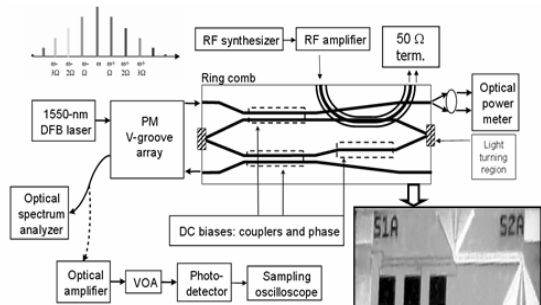


Fig. 1: Ring resonator optical comb schematic view and test setup. Inlay with actual device picture is shown.

Abstract: A novel LiNbO₃ ring resonator cavity is developed for the tunable optical frequency comb generation. Such device enables a high frequency RF modulation of the optical field with independent tuning of the resonator add/drop coupling and phase. The modeling of the integrated ring FCG, fabrication and some experimental results are described.

Introduction

Using an electro-optical phase modulator placed into a resonant cavity is the most direct method of producing a frequency comb that can be also viewed as an optical pulse generator for use, for example, in optical coherent sensing and tomography [1-2]. The modulation frequency should consist of a multiple number of the optical cavity FSR (an integer number of fringe separations) thus efficiently producing a comb by resonantly enhancing the modulation sidebands. High modulation frequency and electro-optical efficiency can provide a relatively broad resulting spectrum. The high bandwidth can be achieved by implementing traveling wave electrodes so that the electrical field co-propagates with the optical field in the LiNbO₃ substrate.

Integrated Ring FCG

In practice, a reliable wide spectral coverage (narrow pulse-train) requires high accuracy stable optical sources along with the well-controlled phase and cavity input/output coupling.

There are some drawbacks limiting the performance of the conventional FCG [3], such as suppression of the total sidebands power caused by increasing the mirror refractivity to increase the relative power transferred to the high order sidebands. By increasing the finesse (and/or modulating index) the pulse-width is narrowed while the peak power remains the same, thus decreasing the sidebands total power (average).

Another difficulty is a small cavity size (1-2cm) (needed for a desired bandwidth) limits the modulating index and overall electro-optical design flexibility. With these issues in mind, a new type of ring-resonator FCG with tunable add/drop coupling was developed [3]. The problem of the high radiative bend losses associated with the weakly-guiding LiNbO₃ ring cavity was addressed by developing relatively low-loss light-reflective regions that helped to overcome the minimal limit for the cavity circumference needed for a reasonably large bandwidth. A schematic view of the ring-cavity OFG and photograph of the actual device are shown Fig. 1.

The analysis of the passive (without RF modulation) ring resonator with the tunable add/drop coupling is described elsewhere [4] for several types of couplers. It was shown that for a realistic lossy cavity the optimal, generally non-identical, add/drop coupling values can be defined and adjusted accordingly via the independent electro-optical perturbations of the couplers and resonator phase. Fig. 2 shows the quality/contrast optimization of such a tunable ring resonator.

3. Ring FCG experiment

The test setup is shown in Fig. 1. Two FCG cavity lengths of ~ 2.25 and ~ 4.5 cm were fabricated, both supporting ~ 6 and ~ 12 GHz modulating frequencies. An external 1550-nm DFB laser source was coupled into one of the bus waveguides of the ring resonator via a V-groove PM fiber array. The thru port of the comb output on the right side was directed to an optical power head using a focusing lens while the reflected port was fiber-coupled via the same V-groove array. RF and DC probes were applied to the traveling-wave, phase, and coupler electrodes of the device. Optical signal from the reflected port was monitored using an optical power meter, optical spectrum analyzer (OSA), and a 30-GHz sampling oscilloscope for

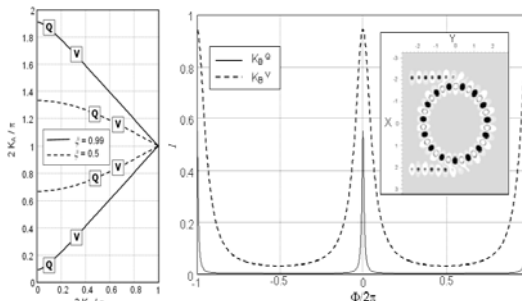


Fig. 2: a) Resonant coupling diagram. Marks indicate the coupling values for the maximum resonator quality Q or contrast V given for two values of cavity losses (1- ξ). b) Q- and V- resonator response optimization by passive design or by an external tuning; for $\xi=0.99$.

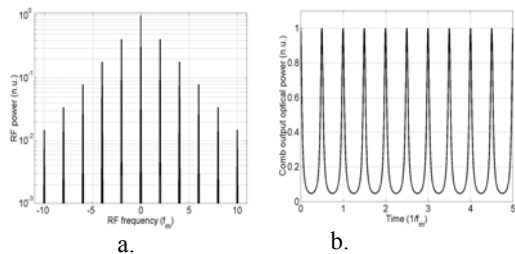


Fig. 6: a) Simulated pulse output of the optical comb biased at resonance. Note the pulse repetition rate is twice the modulation frequency, f_m . b) RF power spectrum of the pulse train shown on the left. Note the absence of odd harmonics of f_m .

Fig. 7 shows a typical plot of the calculated P_o and V_{RF} versus the comb phase bias based on a comb resonator model [2]. The comb phase bias changes the resonant frequencies which affects the resonant condition. In Fig. 7, resonance occurs at comb phase bias of 0 or integer multiples of π . Both P_o and V_{RF} approach to their respective minimum values at resonance. However, P_o and V_{RF} also reach minimum values at off-resonance (e.g., phase= $\pi/2$). Therefore, a resonance control algorithm based on minimizing P_o and V_{RF} may or may not converge to resonance depending on the initial phase bias. In order to ensure convergent to resonance regardless of the initial phase bias, the control algorithm must first adjust the comb bias to maximize P_o .

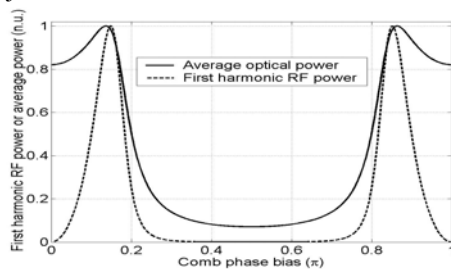


Fig. 7: Average optical power and RF power of the f_m first harmonic or V_{RF} versus phase bias. Resonance corresponds to comb phase bias of zero or integer multiples of π .

Once P_o reaches maximum, the comb bias must be varied to minimize V_{RF} . Note that if the modulation index is high, the relative phase bias position where maximum of P_o and V_{RF} are located may be interchanged. In this case, the algorithm should be modified so that V_{RF} be maximized first followed by minimization of P_o .

Experiments were conducted to verify the operation of the resonant control loop and also validate the performance of the variable ring comb. Fig. 8 shows the pulse waveforms of the ring comb output displayed on the sampling oscilloscope with and without the resonant control loop running. The sampling oscilloscope was set to infinite persistent mode so that drift and fluctuations of the comb output can be accumulated and recorded. Note that the pulse waveforms shown in the figures are limited by the detection bandwidth (40GHz) of the oscilloscope. With the control loop running, one can see from Fig. 8

that the ring comb device is maintained at resonance with a stable output pulse train at a repetition rate of about 24.4 GHz or $2f_m$. On the other hand, without the control loop the comb device seems to fluctuate randomly in and out of resonance causing the output pulse pattern to drift in time. As a result, no stable pulse train at $2f_m$ was observed. The fluctuation can be attributed to variation of the injection laser optical frequency and environmental-induced perturbations on the refractive index of the LN waveguide of the comb device.

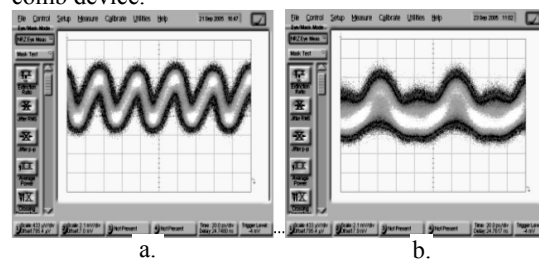


Fig. 8: Pulse train of ring comb output with (a) and without (b) resonant control. Injection laser wavelength: 1544nm. The comb is driven at 12.2GHz. a)- pulse repetition rate is 24.4GHz. Detection bandwidth: 40GHz. Infinite persistent mode \sim 2 minutes. Horizontal scale: 20ps/div.

Conclusions

A novel fully tunable, LiNbO₃-integrated weak-guiding ring FCG device was developed. This device potentially overcomes the limitations of integrated LiNbO₃ Fabry-Perot cavities and existing ring cavities with limited coupling tunability. Preliminary experimental results are consistent with the simulation results. Current and future development includes the add/drop adjustment optimization, cavity losses minimization, and improved resonant control scheme that utilizes both available optical and electrical outputs for a more robust and stable performance.

LiNbO₃-based electro-optical ring resonators with autonomous input/output coupling ratios and phase tuning may become promising components for a variety of applications such as FCG, channel add-drop filters, sensors, chromatic dispersion compensators, optical delay lines, and Er-doped cavities for loss compensation, to mention a few.

Acknowledgments

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